

# Dynamical sensitivity of critical planar percolation, and the Incipient Infinite Cluster

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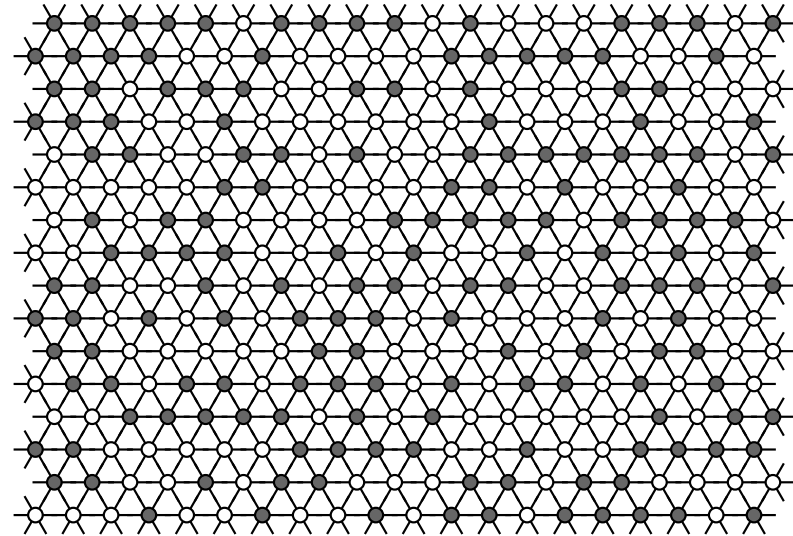
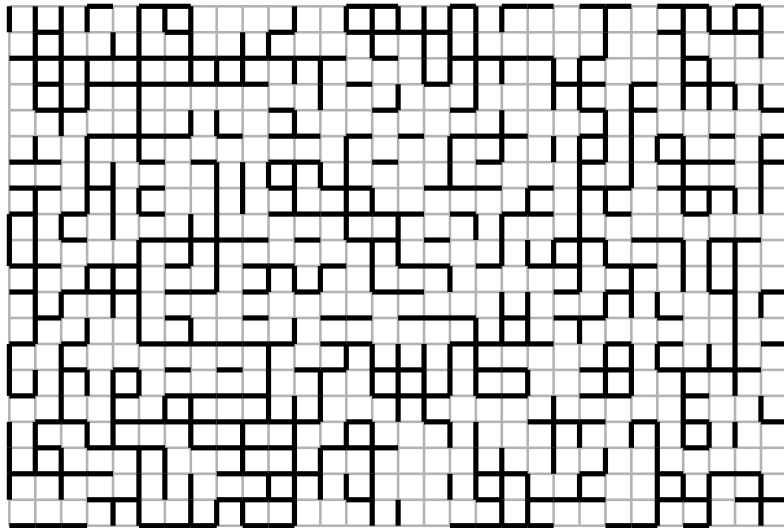
Joint work with **Christophe Garban** (Université Paris-Sud and ENS)  
and **Oded Schramm** (Microsoft Research)  
and **Alan Hammond** (Courant Institute)

## Plan of the talk

- Critical percolation: RSW, conformal invariance,  $SLE_6$  exponents.
- Noise sensitivity of critical percolation.
- Dynamical percolation.
- Why is the Fourier spectrum useful?
- Exceptional times and the IIC.
- The Fourier spectrum of critical percolation. Strategy of proof.
- Further results and questions.

## Bernoulli( $p$ ) site and bond percolation

Given an (infinite) graph  $G = (V, E)$  and  $p \in [0, 1]$ . Each site (or bond) is chosen open with probability  $p$ , closed with  $1 - p$ , independently of each other. Consider the **open connected clusters**.

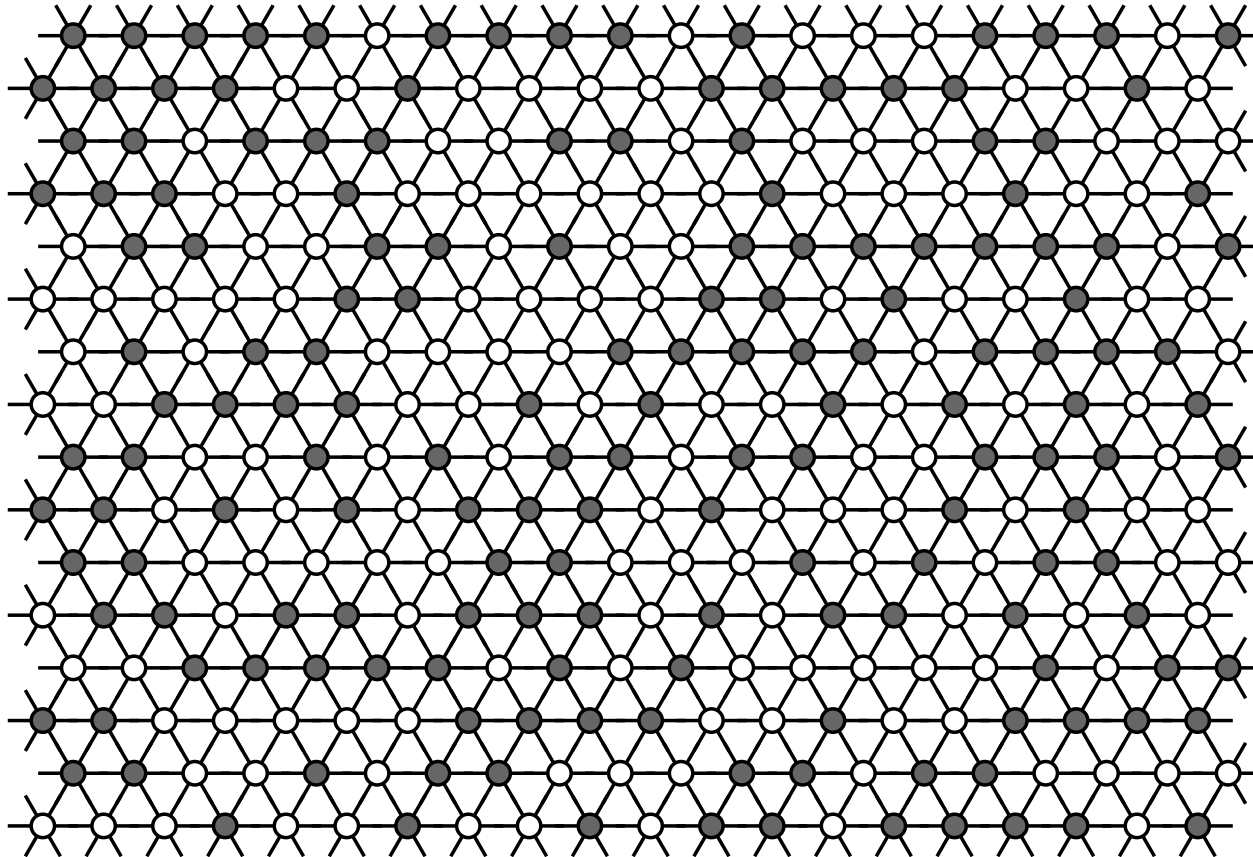


**Theorem (Harris 1960 and Kesten 1980).**

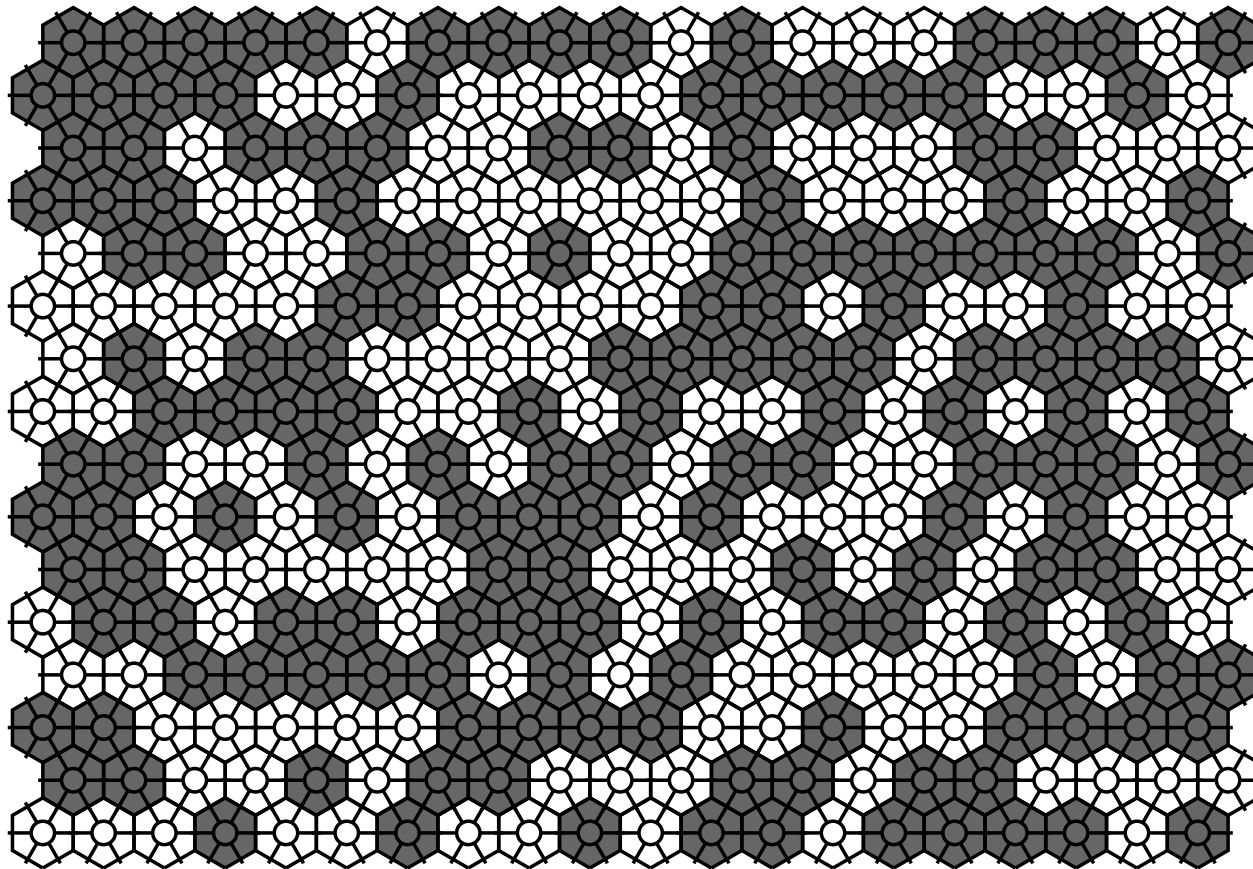
$$p_c(\mathbb{Z}^2, \text{bond}) = p_c(\Delta, \text{site}) = 1/2, \text{ and } \theta(1/2) = 0.$$

For  $p > 1/2$ , there is a.s. one infinite cluster.

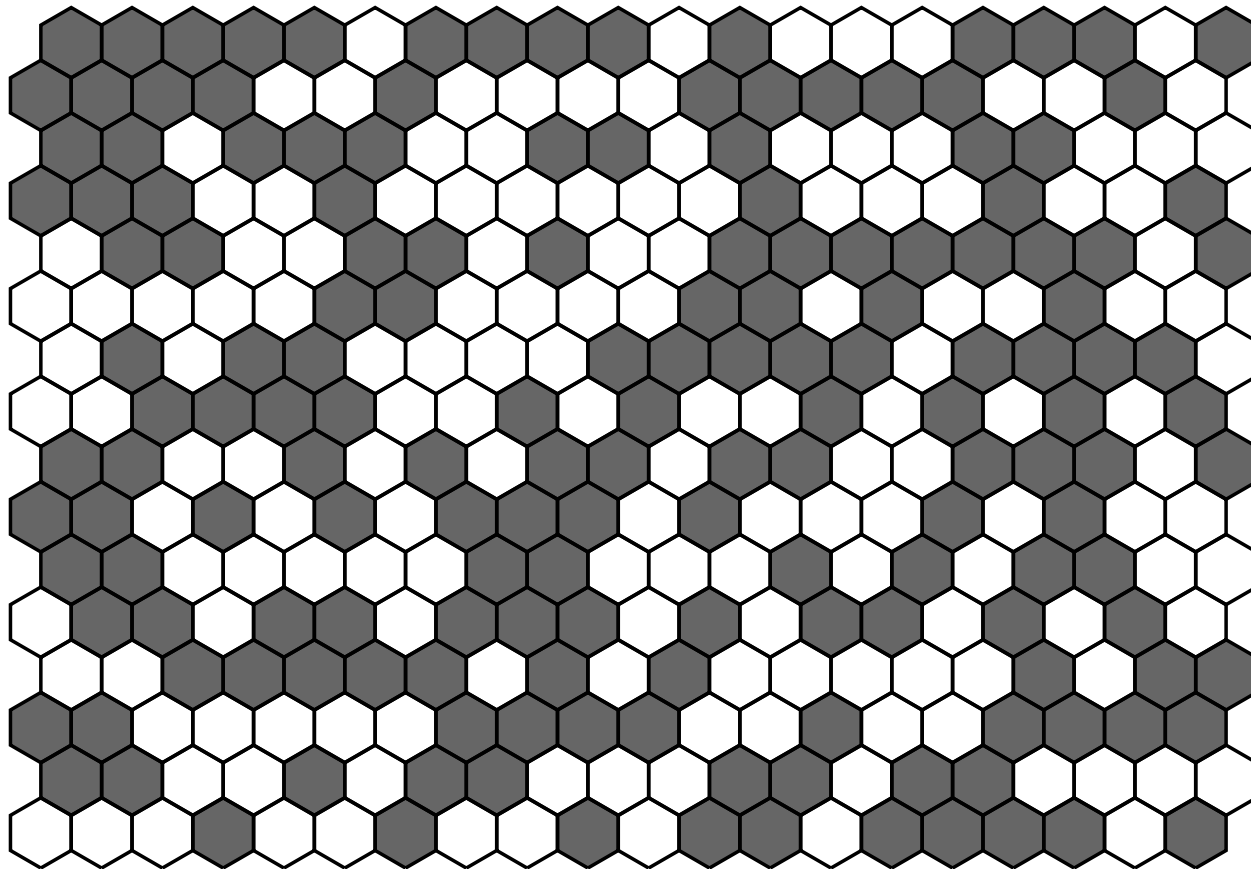
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= face percolation on hexagonal grid:**



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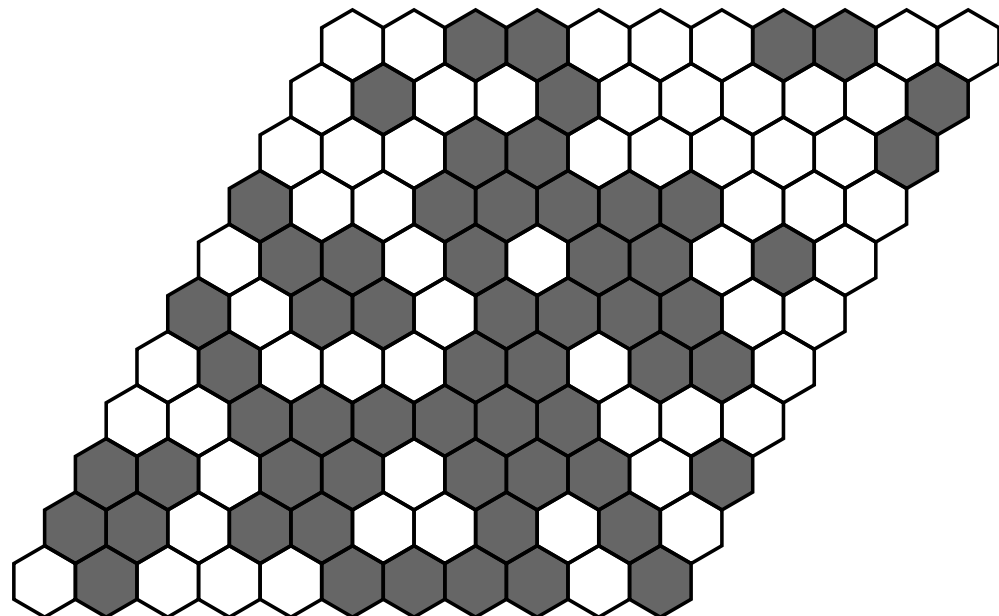
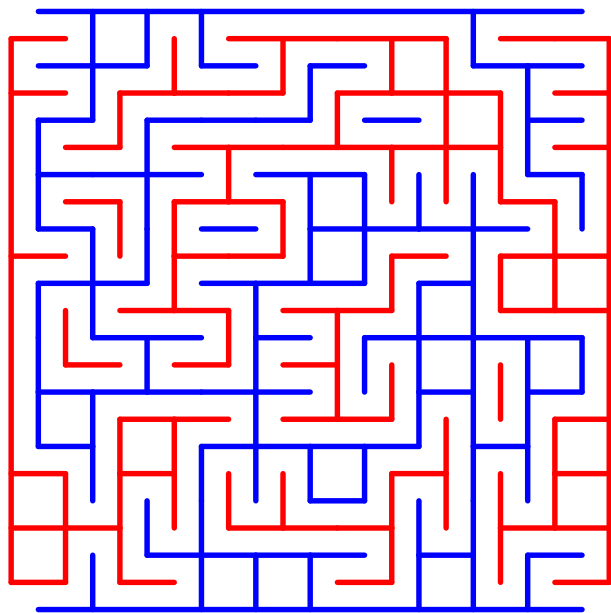
**Site percolation on triangular grid  $\Delta$   
= face percolation on hexagonal grid:**



## Why is $p_c = 1/2$ ? Duality!

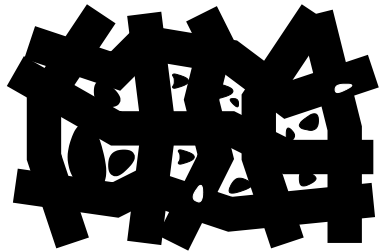
$\mathbb{Z}^2$  bond percolation at  $p = 1/2$ : in an  $(n + 1) \times n$  rectangle, **left-right crossing** has probability exactly  $1/2$ .

For site percolation on  $\Delta$ , same on an  $n \times n$  rhombus.



## Crossing probabilities and criticality

$p \approx 0.8$



$p \approx 0.55$



$p = 0.5$



$p \approx 0.45$



**Theorem (Russo 1978 and Seymour-Welsh 1978).** In critical percolation on any planar lattice, for  $L, n > 0$ ,

$$0 < a_L < \mathbf{P}[\text{left-right crossing in } n \times Ln] < b_L < 1.$$

Same holds for annulus-crossings. It follows easily that

$$(r/R)^\alpha < \mathbf{P}[\partial B_r \longleftrightarrow \partial B_R] < (r/R)^\beta.$$

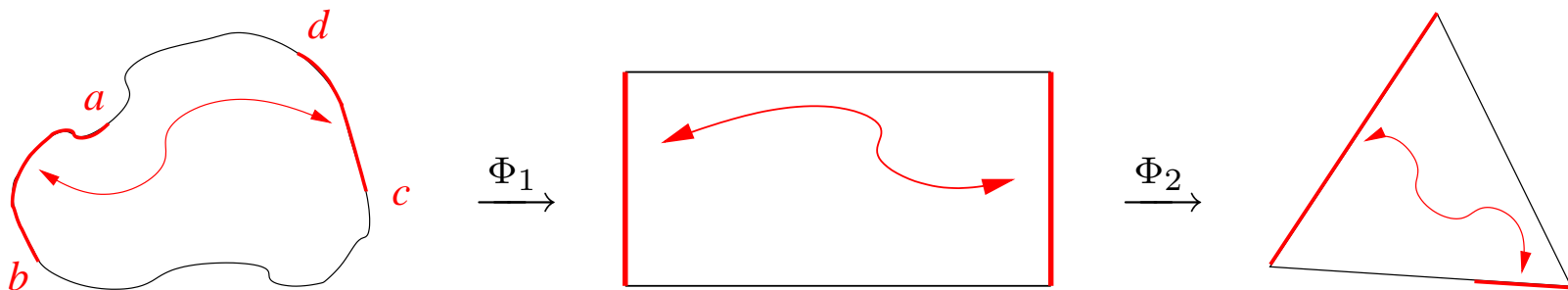


## Conformal invariance on $\Delta$

**Theorem (Smirnov 2001).** For  $p=1/2$  bond percolation on  $\Delta_\epsilon$ , and  $D \subset \mathbb{R}$  simply connected domain with four boundary points  $\{a, b, c, d\}$ ,

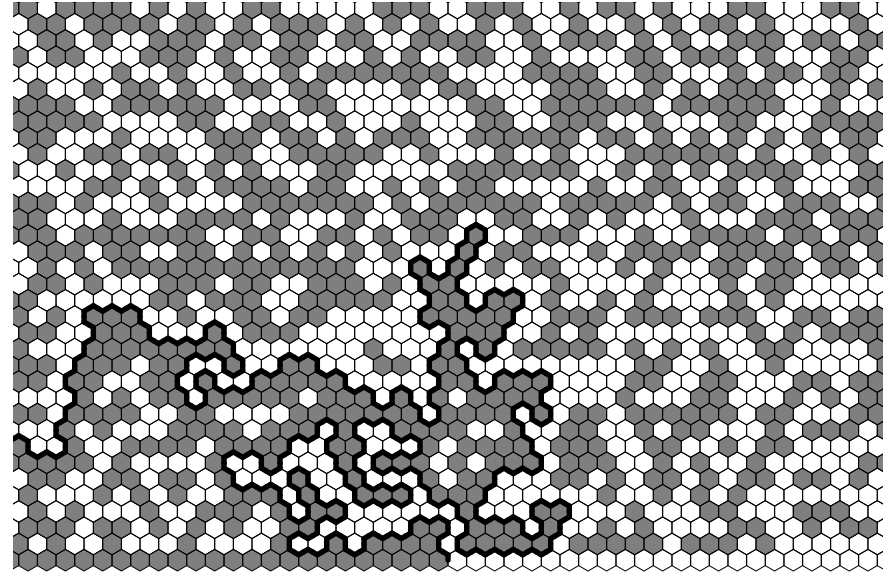
$$\lim_{\epsilon \rightarrow 0} \mathbf{P} \left[ ab \longleftrightarrow cd \text{ inside the discrete approximation } D_\epsilon \right]$$

exists, is strictly between 0 and 1, and is conformally invariant.



## $SLE_6$ exponents

Given the conformal invariance, the exploration path converges to the **Stochastic Loewner Evolution** with  $\kappa = 6$  (Schramm 2000).



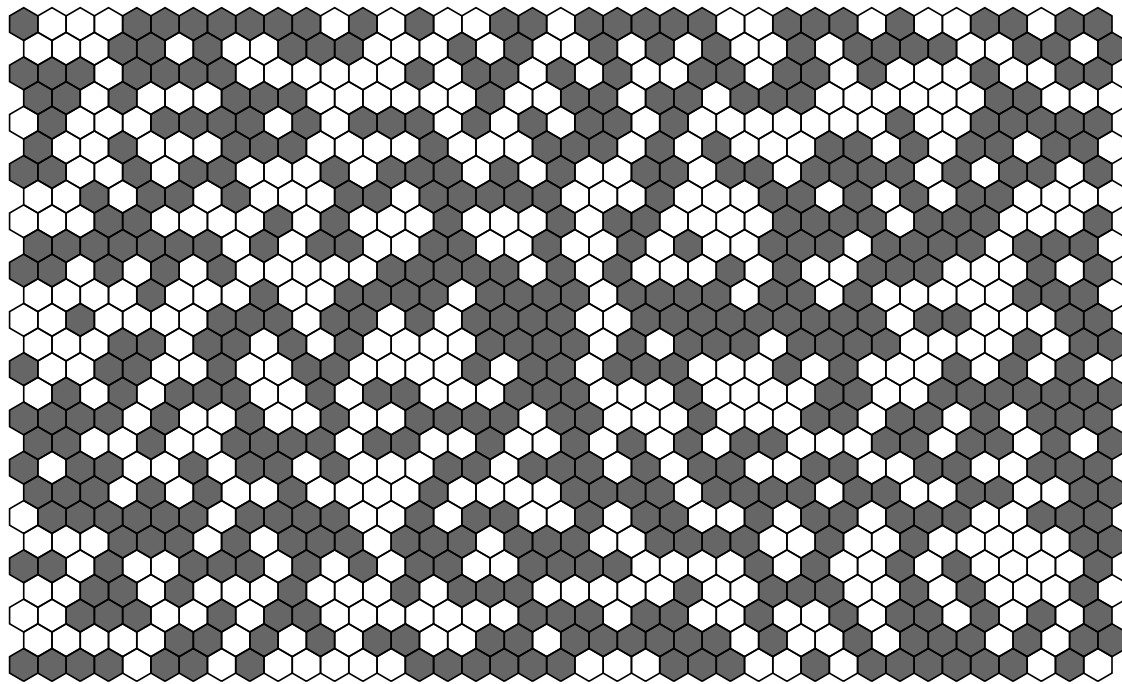
Using the  $SLE_6$  curve, several **critical exponents** can be computed (Lawler-Schramm-Werner, Smirnov-Werner 2001), e.g.:

$$\alpha_4(r, R) := \mathbf{P} \left[ \begin{array}{c} R \\ \text{Diagram of a disk with radius } R \text{ and a smaller disk of radius } r \text{ inside. Two paths, one red and one blue, connect the boundary of the inner disk to the boundary of the outer disk.} \\ r \end{array} \right] = (r/R)^{5/4+o(1)},$$

while  $\alpha_1(r, R) = (r/R)^{5/48+o(1)}$  and  $\mathbf{P}_{p_c+\epsilon}[0 \longleftrightarrow \infty] = \epsilon^{5/36+o(1)}$ .

## Percolation and noise

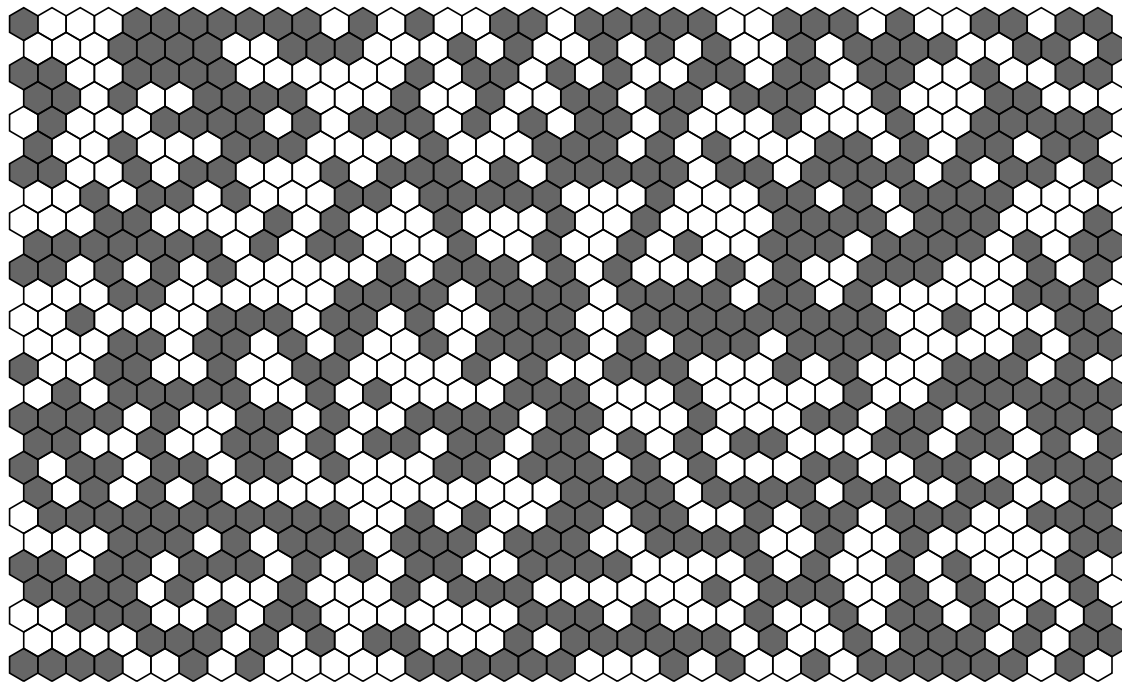
Take an  $\omega$  critical percolation configuration. Let  $\omega^\epsilon$  be a new configuration, where each site (or bond) is **resampled** with probability  $\epsilon$ , independently. (The  $\epsilon$ -noised version of  $\omega$ .)



For how large an  $\epsilon$  can we still predict from  $\omega$  whether there is a left-right crossing in  $\omega^\epsilon$ ?

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## Noise sensitivity of percolation

**Theorem (Benjamini, Kalai & Schramm 1998).** If  $\epsilon > 0$  is fixed, and  $f_n$  is the indicator function for a left-right percolation crossing in an  $n \times n$  square, then as  $n \rightarrow \infty$

$$\mathbf{E} [ f_n(\omega) f_n(\omega^\epsilon) ] - \mathbf{E} [ f_n(\omega) ]^2 \rightarrow 0.$$

This holds for all  $\epsilon = \epsilon_n > c/\log n$ .

**Theorem (Steif & Schramm 2005).** Same if  $\epsilon_n > n^{-a}$  for some positive  $a > 0$ . If triangular lattice, may take any  $a < 1/8$ .

**Theorem (Garban, P & Schramm 2008).** Same holds if and only if  $\epsilon_n \mathbf{E} [ |\text{pivotal}| ] \rightarrow \infty$ . For triangular lattice, this threshold is  $\epsilon_n = n^{-3/4+o(1)}$ .

## Naive idea: how many pivotals are there?

A site (or bond) is **pivotal** in  $\omega$ , if flipping it changes the existence of a left-right crossing. Expected number:  $\mathbf{E}|\text{Piv}_n| \asymp n^2 \alpha_4(n) \quad (= n^{3/4+o(1)})$ .

Second moment:  $\mathbf{E}[|\text{Piv}_n|^2] \leq C (\mathbf{E}|\text{Piv}_n|)^2$ .

From Chebyshev:  $\mathbf{P}[|\text{Piv}_n| > \lambda \mathbf{E}|\text{Piv}_n|] < C/\lambda^2$  for all  $\lambda > 0$ -ra.

From Cauchy-Schwarz:  $\exists \delta > 0$ , s.t.  $\mathbf{P}[|\text{Piv}_n| > \delta \mathbf{E}|\text{Piv}_n|] > \delta$ .

Moreover:  $\mathbf{P}[0 < |\text{Piv}_n| < \lambda \mathbf{E}|\text{Piv}_n|] \asymp \lambda^{11/9+o(1)}$ , if  $\lambda < 1$  (on  $\Delta$ ).

Cannot have many pivotals.  $\implies$  If  $\epsilon_n \mathbf{E}[|\text{Piv}_n|] \rightarrow 0$ , then we don't hit any pivotals.  $\implies$  Asymptotically full correlation.

Cannot have few pivotals.  $\implies$  If  $\epsilon_n \mathbf{E}[|\text{Piv}_n|] \rightarrow \infty$ , then we do hit many pivotals. But this  $\not\implies$  asymptotic independence!

# Dynamical percolation

Each variable is resampled according to an independent Poisson(1) clock. This is a Markov process  $\{\omega(t) : t \in [0, \infty)\}$ , in which  $\omega(t + s)$  is an  $\epsilon$ -noised version of  $\omega(t)$ , with  $\epsilon = 1 - \exp(-s)$ .

An **exceptional time** is such a (random)  $t$ , at which an almost sure property of the static process fails for  $\omega(t)$ .

**Main example:** (Non-)existence of an infinite cluster in percolation.

**Toy example:** Brownian motion on the circle does sometimes hit a fix point, as opposed to its static version: a uniform random point.

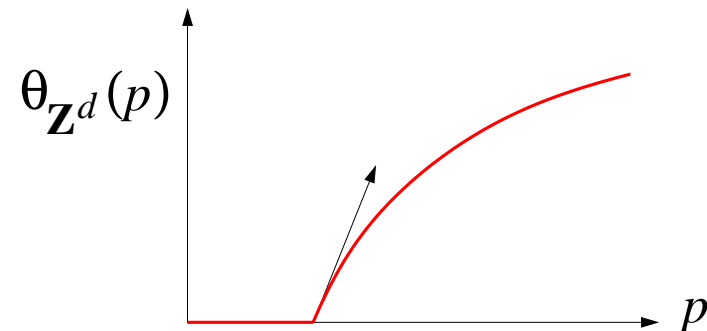
In this toy example, the set of exceptional times is a random Cantor set of Lebesgue measure zero (because of Fubini) and Hausdorff-dimension  $1/2$ .

# Dynamical percolation results

Theorem (**Häggström, Peres & Steif 1997**).

- No exceptional times when  $p \neq p_c$ .
- No exceptional times when  $p = p_c$  for bond percolation on  $\mathbb{Z}^d$ ,  $d \geq 19$ .

The second fact is essentially due to:





**Theorem (Steif & Schramm 2005).**

- There are exceptional times (a.s.) for critical site percolation on the triangular lattice.
- They have Hausdorff dimension in  $[1/6, 31/36]$ .

**Theorem (Garban, P & Schramm 2008).**

- There are exceptional times also on  $\mathbb{Z}^2$ .
- On the triangular grid they have Hausdorff dimension  $31/36$ .
- On the triangular grid, there are exceptional times with an infinite white **and** an infinite black cluster simultaneously. (Dim=  $2/3$ .)

## What is the Fourier spectrum and why is it useful?

$f_n : \{\pm 1\}^{V_n} \leftrightarrow \{\pm 1\}$  indicator function of left-right crossing.  $\mathbf{E}[f_n] = 0$ .

$(N_\epsilon f)(\omega) := \mathbf{E}[f(\omega^\epsilon) | \omega]$  is the **noise operator**, acting on the space  $L^2(\Omega, \mu)$ , where  $\Omega = \{\pm 1\}^{V_n}$ ,  $\mu$  uniform measure, inner product  $\mathbf{E}[fg]$ .

Correlation:  $\mathbf{E}[f(\omega^\epsilon)f(\omega)] - \mathbf{E}[f(\omega)]\mathbf{E}[f(\omega^\epsilon)] = \mathbf{E}[f(\omega)N_\epsilon f(\omega)] - \mathbf{E}[f(\omega)]^2$ . So, we would like to **diagonalize** the noise operator  $N_\epsilon$ .

Let  $\chi_v$  be the function  $\chi_v(\omega) = \omega(v)$ ,  $\omega \in \Omega$ .

For  $S \subset V_n$ , let  $\chi_S := \prod_{v \in S} \chi_v$ , the **parity inside  $S$** . Then

$$N_\epsilon \chi_v = (1 - \epsilon) \chi_v; \quad N_\epsilon \chi_S = (1 - \epsilon)^{|S|} \chi_S.$$

Moreover, the family  $\{\chi_S, S \subset V\}$  is an **orthonormal basis** of  $L^2(\Omega, \mu)$ .

Any function  $f \in L^2(\Omega, \mu)$  in this basis (**Fourier-Walsh series**):

$$\hat{f}(S) := \mathbf{E}[f\chi_S]; \quad f = \sum_{S \subset V} \hat{f}(S) \chi_S.$$

The correlation:

$$\begin{aligned} \mathbf{E}[fN_\epsilon f] - \mathbf{E}[f]^2 &= \sum_S \sum_{S'} \hat{f}(S) \hat{f}(S') \mathbf{E}[\chi_S N_\epsilon \chi_{S'}] - \mathbf{E}[f\chi_\emptyset]^2 \\ &= \sum_{\emptyset \neq S \subset V} \hat{f}(S)^2 (1 - \epsilon)^{|S|} = \sum_{k=1}^{|V_n|} (1 - \epsilon)^k \sum_{|S|=k} \hat{f}(S)^2. \end{aligned}$$

May consider the associated “**spectral measure**”  $\nu_f(S) := \hat{f}(S)^2$ . For the crossing function,  $\mathbf{E}[f_n^2] = 1$ , so Parseval says this is a probability measure. A random sample  $\mathcal{S}_n = \mathcal{S}(f_n) \subset V_n$  is a strange random set of bits.

If  $\nu[|\mathcal{S}_n| < tk_n] \rightarrow 0$  as  $t \rightarrow 0$ , then  $(1 - \epsilon)^k \sim \exp(-\epsilon k)$  implies that for  $\epsilon_n \gg 1/k_n$  we have  $\mathbf{E}[f_n N_\epsilon f_n] - \mathbf{E}[f_n]^2 \rightarrow 0$ , **asymptotic independence**.

## Proving existence of exceptional times

### Second Moment Method:

Let  $Q_R := \{t \in [0, 1] : 0 \longleftrightarrow_t R\}$  and  $Z_R := \text{Leb}(Q_R)$ .

$$\mathbf{P}[Q_R \neq \emptyset] = \mathbf{P}[Z_R > 0] \geq \frac{\mathbf{E}[Z_R]^2}{\mathbf{E}[Z_R^2]}.$$

$$\mathbf{E}[Z_R] = \int_0^1 \mathbf{P}[0 \longleftrightarrow_t R] dt = \mathbf{P}[0 \longleftrightarrow R].$$

$$\mathbf{E}[Z_R^2] = \int_0^1 \int_0^1 \mathbf{P}[0 \longleftrightarrow_t R, 0 \longleftrightarrow_s R] ds dt \asymp \int_0^1 \mathbf{E}[f_R(\omega_0) f_R(\omega_s)] ds.$$

Thus we again want to estimate the correlation  $\mathbf{E}[f_R(\omega_0) f_R(\omega_s)] = \mathbf{E}[f_R T_s f_R]$  from above, where

$$T_s f(\omega) := \mathbf{E}[f(\omega_s) \mid \omega_0 = \omega] = N_{1-\exp(-s)} f(\omega).$$

## Hausdorff dimension of exceptional times

If  $f_R$  is the 0/1 indicator of the  $\ell$ -arm event to radius  $R$ , with exponent  $\mathbf{E}[f_R] = R^{-\xi_\ell + o(1)}$ , then **[GPS]**:

$$\mathbf{E}[f_R(\omega_s) f_R(\omega_t)] / \mathbf{E}[f_R(\omega)]^2 \leq O(1) |t - s|^{-(4/3)\xi_\ell + o(1)},$$

as  $|t - s| \rightarrow 0$ . Now, by the **Mass Distribution Principle**, if

$$\sup_R \int_0^1 \int_0^1 \frac{\mathbf{E}[f_R(\omega_t) f_R(\omega_s)]}{\mathbf{E}[f_R(\omega)]^2 |t - s|^\gamma} dt ds < \infty,$$

then  $\dim(\mathcal{E}_\ell) \geq \gamma$  a.s. Hence  $\dim(\mathcal{E}_\ell) \geq 1 - 4\xi_\ell/3$ .

For  $\mathbb{Z}^2$ , we have “ $\xi_1 + \xi_4 < \xi_5 = 2$ ”, hence  $1 - \frac{\xi_1}{2 - \xi_4} > 0$ , so there exist exceptional times.

## Local time measure for exceptional times [HPS]

$$\overline{M}_r(\omega_s) := \frac{\mathbb{1}\{0 \leftrightarrow r\}}{\mathbf{P}[0 \leftrightarrow r]}, \quad \overline{\mu}_r[a, b] := \int_a^b \overline{M}_r(\omega_s) ds, \quad \overline{\mu}[a, b] := \lim_{r \rightarrow \infty} \overline{\mu}_r[a, b].$$

This  $\overline{M}_r(\omega)$  is a martingale w.r.t. the filtration  $\overline{\mathcal{F}}_r$  of the percolation space generated by the variables  $\mathbb{1}\{0 \leftrightarrow r\}$ . Moreover,  $\mathbf{E}\overline{\mu}_r[a, b] = b - a$ , and, by the 2nd Moment Estimate,  $\sup_r \mathbf{E}[\overline{\mu}_r[a, b]^2] < C_1$ . So  $\lim_r$  exists.

$$M_H(\omega) := \lim_{R \rightarrow \infty} \frac{\mathbf{P}[0 \leftrightarrow R \mid \omega^H]}{\mathbf{P}[0 \leftrightarrow R]} = \lim_{R \rightarrow \infty} \frac{\mathbf{P}[\omega^H \mid 0 \leftrightarrow R]}{\mathbf{P}[\omega^H]} = \frac{\text{IIC}(\omega^H)}{\mathbf{P}[\omega^H]},$$

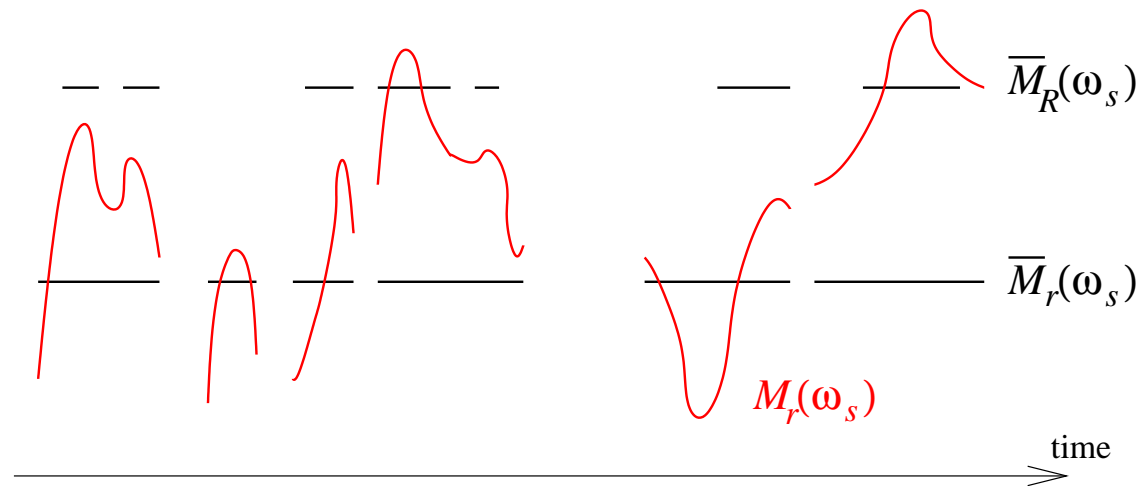
where  $H \subset \Delta$ , and IIC is Kesten's **Incipient Infinite Cluster** measure (1986).

$$M_r(\omega_s) := M_{B_r}(\omega_s), \quad \mu_r[a, b] := \int_a^b M_r(\omega_s) ds, \quad \mu[a, b] := \lim_{r \rightarrow \infty} \mu_r[a, b].$$

Now  $M_r(\omega)$  is a MG w.r.t. the full filtration  $\mathcal{F}_r$  generated by  $\omega(B_r)$ , again  $\mathbf{E}\mu_r[a, b] = b - a$ , and  $M_r(\omega) \leq C_2 \overline{M}_r(\omega)$  because of quasi-multiplicativity:

$$\begin{aligned} \frac{\mathbf{P}[0 \leftrightarrow R | \omega^{B_r}]}{\mathbf{P}[0 \leftrightarrow R]} &\asymp \frac{\mathbf{P}[0 \leftrightarrow R | \omega^{B_r}]}{\mathbf{P}[0 \leftrightarrow r] \mathbf{P}[r \leftrightarrow R]} \\ &\leq \frac{\mathbf{P}[r \leftrightarrow R | \omega^{B_r}] \mathbb{1}\{0 \leftrightarrow r\}}{\mathbf{P}[0 \leftrightarrow r] \mathbf{P}[r \leftrightarrow R]} = \frac{\mathbb{1}\{0 \leftrightarrow r\}}{\mathbf{P}[0 \leftrightarrow r]}. \end{aligned}$$

Hence, both local time measures exist, and are clearly supported inside  $\mathcal{E}$ .



**Theorem (Hammond, P & Schramm 2008).** The two measures coincide:  $\bar{\mu} = \mu$  a.s. At a  $\mu$ -typical time, the configuration has the distribution of IIC. But the configuration at the first exceptional time is different.

Question: is it true that  $\text{supp}(\mu) = \mathcal{E}$ ?

## Basic properties of the spectral sample

**Inclusion formula:**  $\nu_f[\mathcal{S} \subset U] = \mathbf{E}\left[\mathbf{E}[f | U]^2\right]$ .

Proof:

$$\mathbf{E}[\chi_S | U] = \begin{cases} \chi_S & S \subset U, \\ 0 & S \not\subset U. \end{cases}$$

Thus  $\mathbf{E}\left[\mathbf{E}[f | U]^2\right] = \mathbf{E}\left[\left(\sum_{S \subset U} \hat{f}(S) \chi_S\right)^2\right] = \sum_{S \subset U} \hat{f}(S)^2$ . ■

From this, for disjoint subsets  $A$  and  $B$ ,

$$\begin{aligned} \nu[\mathcal{S} \cap B \neq \emptyset, \mathcal{S} \cap A = \emptyset] &= \nu[\mathcal{S} \subseteq A^c] - \nu[\mathcal{S} \subseteq (A \cup B)^c] \\ &= \mathbf{E}\left[\mathbf{E}[f | A^c]^2 - \mathbf{E}[f | (A \cup B)^c]^2\right] \\ &= \mathbf{E}\left[\left(\mathbf{E}[f | A^c] - \mathbf{E}[f | (A \cup B)^c]\right)^2\right]. \end{aligned}$$



## For the spectral sample $\mathcal{S}_L$ of the $L \times L$ crossing:

With  $A := \emptyset$  we get:  $\nu[\mathcal{S}_L \cap B \neq \emptyset] \leq C \alpha_4(B, V_L)$ ;

with  $A := B^c$  we get:  $\nu[\emptyset \neq \mathcal{S}_L \subseteq B] \leq C \alpha_4(B, V_L)^2$ .

If  $B = \{x\}$ : equality in both cases. Hence  $\nu[x \in \mathcal{S}_L] = \mathbf{P}[x \in \text{Piv}_L]$ , and

$$\mathbf{E}_\nu[|\mathcal{S}_L|] = \mathbf{E}[|\text{Piv}_L|] =: m_L \quad (= L^{3/4+o(1)}).$$

If  $B$  is a sub-square of side  $r$ , and  $B' = B/3$ , then

$$\begin{aligned} \mathbf{E}_\nu \left[ |\mathcal{S} \cap B'| \mid \mathcal{S} \cap B \neq \emptyset \right] &= \sum_{x \in B'} \frac{\nu[x \in \mathcal{S}]}{\nu[\mathcal{S} \cap B \neq \emptyset]} \geq \sum_{x \in B'} \frac{4\alpha_4(x, V_L)}{\alpha_4(B, V_L)} \\ &\asymp \sum_{x \in B'} \alpha_4(x, B) \asymp r^2 \alpha_4(r) \asymp m_r, \end{aligned}$$

as we would expect from a random fractal-like set. But we need something stronger: with good probability, and conditioned on other sub-squares.

## Main results for the spectral sample (GPS)

If  $r \in [1, L]$ , then  $\{|\mathcal{S}_L| < m_r\}$  is basically equivalent to being contained inside some  $r \times r$  sub-square:

$$\mathbf{P} [ |\mathcal{S}_L| < m_r ] \asymp \alpha_4(r, L)^2 \left( \frac{L}{r} \right)^2.$$

In particular, on the triangular lattice  $\Delta$ ,

$$\mathbf{P} [ |\mathcal{S}_L| < \lambda m_L ] \asymp \lambda^{2/3+o(1)}.$$

The *scaling limit* of  $\mathcal{S}_L$  is a conformally invariant Cantor-set with Hausdorff-dimension  $3/4$ .

The existence of the scaling limit follows from **Schramm & Smirnov**: *Percolation is black noise*, answering a question of Tsirelson.

## The strategy of proof

Tile the  $L \times L$  square with  $(L/r)^2$  boxes of size  $r$ . Let  $X = X_{r,L}$  be the number of boxes intersecting  $\mathcal{S}_L$ . We already know that

$$\mathbf{E}[X] \geq \alpha_4(r, L)(L/r)^2 \asymp (L/r)^{3/4+o(1)}.$$

**1st step:**  $X$  is smaller than  $C \log(L/r)$  with only very small probability.

**2nd step:** In a non-empty  $r$ -box, with positive probability  $|\mathcal{S}_L| \geq c m_r$ .

If we could repeat this step for each of the  $X$  nonempty boxes,  $\mathcal{S}_L$  would be large almost surely.

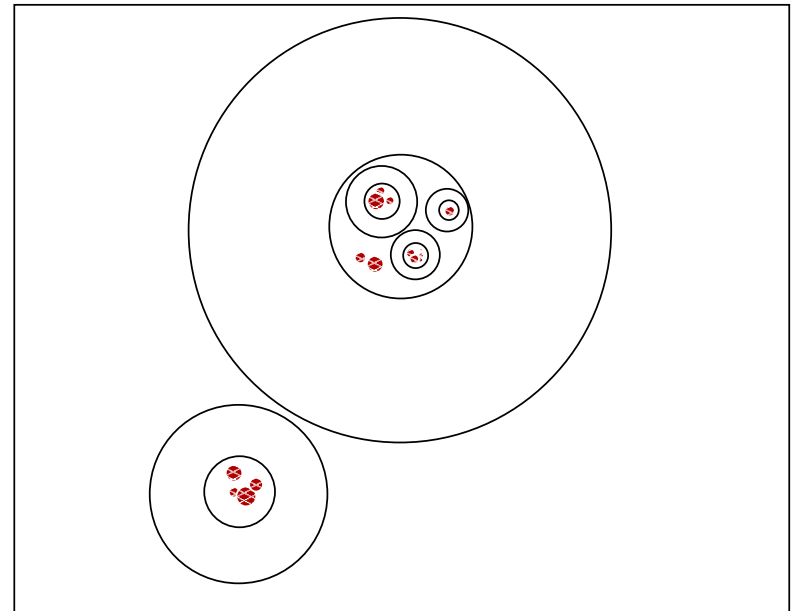
But we can prove Step 2 only in the presence of negative information about  $\mathcal{S}_L$  everywhere else! (Partial independence.)

**3rd step:** Using a sampling trick and a strange large deviation result, 1+2 turns out to be enough.

## Annulus structures

**Proposition 1.**  $\nu[X \leq k] \leq k^{C \log k} (L/r)^2 \alpha_4(r, L)^2.$

An annulus structure  $\mathcal{A}$  compatible with a set  $S$ :

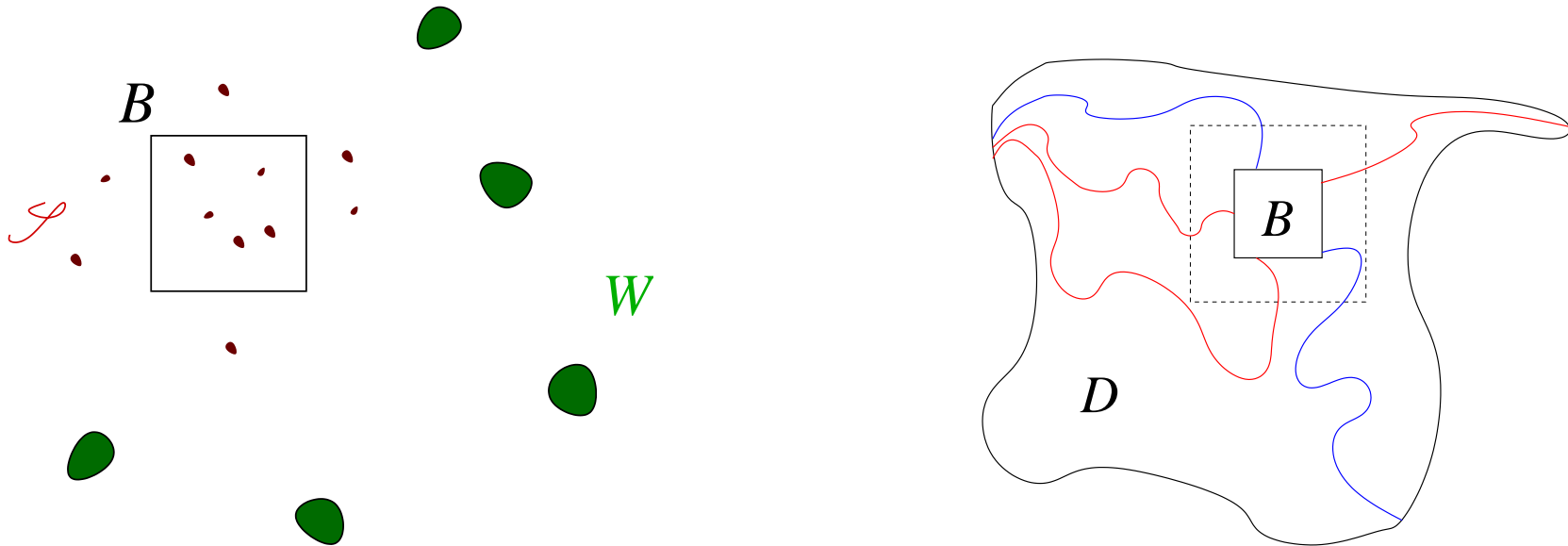


**Lemma.**  $\nu[\mathcal{S} \text{ compatible with } \mathcal{A}] \leq \prod_{A \in \mathcal{A}} \alpha_4(A)^2.$

Thus, we need to construct a family of annulus structures that has some member compatible with any  $k$ -element set, but  $\sum_{\mathcal{A}} \prod_{A \in \mathcal{A}} \alpha_4(A)^2$  is still small. This is done recursively.

## Partial independence

**Proposition 2.** If  $B$  is an  $r$ -box in  $[0, L]^2$ , and  $W \cap B = \emptyset$ , then  $\mathbf{P}\left[|\mathcal{S} \cap B| > cr^2 \alpha_4(r) \mid \mathcal{S} \cap W = \emptyset \neq \mathcal{S} \cap B\right] \geq c$ .



**Separation Lemma.** If  $\text{dist}(B, \partial D) > \text{diam}(B)$ , then conditioned on the  $k$ -arm event in  $D \setminus B$  with fixed endpoints on  $\partial D$ , then with a uniformly positive conditional probability the  $k$  arms are “well-separated” around  $B$ .

## Large deviation lemma

**Proposition 3.** Suppose  $X_i, Y_i \in \{0, 1\}$ ,  $i = 1, \dots, n$ , and that  $\forall J \subset [n]$  and  $\forall i \in [n] \setminus J$

$$\mathbf{P}[Y_i = 1 \mid \forall_{j \in J} Y_j = 0] \geq c \mathbf{P}[X_i = 1 \mid \forall_{j \in J} Y_j = 0].$$

Then

$$\mathbf{P}[\forall_i Y_i = 0] \leq c^{-1} \mathbf{E} \left[ \exp \left( -(c/e) \sum_i X_i \right) \right].$$

We use this with  $X_j := 1_{\mathcal{S} \cap B_j \neq \emptyset}$  and  $Y_j := 1_{\mathcal{S} \cap B_j \cap Q \neq \emptyset}$  for a random Bernoulli set  $Q$ , independent from everything else, with density so that it meets with probability  $1/2$  a fixed set of cardinality  $m_r$  in  $B_j$ .

## Some related results and questions

**Theorem (Garban, P & Schramm 2008).** On  $\Delta_\eta$ , with rate  $\eta^{3/4+o(1)}$  clocks, the **scaling limit of the dynamical percolation** exists as a Markov process, and can be understood via the pivotals at the visible scales. It is *conformally covariant*: if the domain is changed by  $\phi(z)$ , then time is scaled locally by  $|\phi'(z)|^{3/4}$ .

Also, the scaling limit of **near-critical percolation** exists, and is conformally covariant. Consequently, the scaling limit of the **Minimal Spanning Tree** exists and is rotationally and scale invariant, but *not* conformally.

The Fourier work gives that in the scaling limit, the correlation of left-right crossing between times 0 and  $t$  is  $t^{-2/3+o(1)}$ . Is the probability of having the crossing all the way between time 0 and  $t$  is  $\exp(-t^{2/3+o(1)})$ ?

Similar proofs for other Boolean functions?

Crossing function, but non-unif measure, e.g. **Random Cluster models**? Ising is expected to be stable, because of non-existence of pivotals ( $\kappa < 4$ ).