

10.1 *DLR condition and conditional probability.* Let  $\mu_{\Lambda}^{\beta, \beta'}(\cdot | \eta)$  denote the grand canonical measure of some system in the box  $\Lambda$  with the boundary condition  $\eta \in \Omega_{\Lambda^c}$ . Using this, define

$$\gamma_{\Lambda}^{\beta, \beta'}(\cdot | \eta) := \mu_{\Lambda}^{\beta, \beta'}(\cdot | \eta \cap \Lambda^c) \otimes \delta_{\eta \cap \Lambda^c}$$

for all  $\eta \in \Omega$  as a product measure on  $\Omega = \Omega_{\Lambda} \times \Omega_{\Lambda^c}$ .

The Dobrushin-Lanford-Ruelle condition for a measure  $\mu$  on  $\Omega$  to be Gibbs is

$$\mu = \mu \otimes \gamma_{\Lambda}^{\beta, \beta'} \quad \text{for every bounded } \Lambda,$$

where  $\gamma_{\Lambda}^{\beta, \beta'}$  is viewed as a probability kernel from  $\Omega$  to  $\Omega$ . Show that this is the same as requiring that

$$\mu_1 = \mu_2 \otimes \mu_{\Lambda}^{\beta, \beta'},$$

where  $\mu_1$  and  $\mu_2$  are the two marginals of the measure  $\mu$  on  $\Omega = \Omega_{\Lambda} \times \Omega_{\Lambda^c}$ , and  $\mu_{\Lambda}^{\beta, \beta'}$  is viewed as a probability kernel from  $\Omega_{\Lambda^c}$  to  $\Omega_{\Lambda}$ .

10.2 **(homework)** *The partition function and the effect of ignoring the velocity.* Consider a system of interacting point particles in  $d$  dimensions with Hamiltonian  $H(q, p) = \sum_i \frac{\vec{p}_i^2}{2m} + \sum_{i < j} \Phi(|q_i - q_j|)$ . Consider the canonical partition function  $Z^{can}(V, N, \beta)$  and the grand canonical partition function  $Z^{gr}(V, \beta, \beta')$  (or  $Z^{gr}(V, \beta, z)$  if you like), which are integrals of some weight functions on the phase space.

Now define the *configurational partition functions* by “ignoring the velocity”. That is, e.g.  $Z_{conf}^{can}(V, N, \beta)$  is an integral on the *configuration space only* of the canonical weight function with the kinetic energy omitted.

- (a) Find the relation between the partition function and the configurational partition function, both in the canonical and the grand canonical setting.
- (b) Calculate the density of the free gas with parameters  $\beta$ ,  $z$  and  $m$ .

10.3 **(homework)** *Is the grand canonical ensemble with boundary condition well defined?* Show that the grand canonical measure is well defined with any boundary condition in a system of interacting point particles with a bounded finite range pair interaction. What can we say if the pair interaction is only tempered and stable?

10.4 *Consistency property.* Show that  $\gamma_{\Lambda'}^{\beta, \beta'} = \gamma_{\Lambda}^{\beta, \beta'} \otimes \gamma_{\Lambda'}^{\beta, \beta'}$  for every bounded  $\Lambda \subset \Lambda' \subset \mathbb{R}^d$ .

10.5 *Gibbs measures of the free gas.* Find and describe every Gibbs measure of the free gas.

10.6 **(homework)** *Ising model in one dimension.* For the Ising model in one dimension, let the phase space be  $\omega = \{-1, 1\}^N$  and the Hamiltonian be  $H : \Omega \rightarrow \mathbb{R}$  be defined as

$$H(\sigma_1, \dots, \sigma_N) := \sum_{i=1}^N (J\sigma_i\sigma_{i+1} + h\sigma_i).$$

(We use the convention  $\sigma_{N+1} := \sigma_1$ , which corresponds to periodic boundary conditions.)

- (a) Calculate the partition function

$$Z(N, \beta, h) := \sum_{\sigma \in \Omega} e^{-\beta H(\sigma)}.$$

*Hint:* In the expression defining  $Z$ , discover a power of a  $2 \times 2$  matrix. If you do it well, this matrix will be symmetric. In the end, you only need to calculate the eigenvalues.

- (b) Calculate the entropy in the thermodynamic limit (understand the question well) and show that it is a smooth function of  $h$  for every  $\beta > 0$ .